MATHEMATICAL MODELLING AND ANALYSIS Volume 26, Issue 2, 337–338, 2021 https://doi.org/10.3846/mma.2021.14527



Corrigendum Correction to the Paper: An Energy Dissipative Spatial Discretization for the Regularized Compressible Navier-Stokes-Cahn-Hilliard System of Equations (in *Math. Model. Anal.*, 25(1): 110–129, https://doi.org/10.3846/mma.2020.10577)

Vladislav Balashov^a and Alexander Zlotnik^{b,a}

^a Keldysh Institute of Applied Mathematics Miusskaya sqr., 4, 125047 Moscow, Russia
^b Higher School of Economics University Pokrovskii bd. 11, 109028 Moscow, Russia E-mail(corresp.): azlotnik@hse.ru E-mail: vladislav.balashov@gmail.com

Received March 2, 2021; accepted March 25, 2021

Abstract. We correct the proof of Theorem 2 in the mentioned paper concerning finite-difference equilibrium solutions.

Keywords: regularized viscous compressible Navier-Stokes-Cahn-Hilliard equations, finitedifference discretization in space, equilibrium solutions.

AMS Subject Classification: 65M06,35Q35.

In this note, we correct the proof of Theorem 2 (p. 120–121) in [1] following our recent paper [2]. Below we exploit the notation from [1].

We consider the equilibrium solutions $(\rho, \mathbf{u}, C) = (\rho_S, \mathbf{u}_S, C_S) \in H_X(\omega_h)$ with $\rho_S > 0$, $\mathbf{u}_S = 0$ and $0 < C_S(x) < 1$. For them, the following equations hold

$$\delta_i^* J_{iS} = 0, \tag{0.1}$$

$$s_l^* \{ (s_l \rho) [\delta_l (G_h - \Phi) - (s_l \mu) \delta_l C] \} = 0, \ l = 1, 2, 3, \tag{0.2}$$

$$\delta_i^* \left[J_{iS} s_i C - M(s_i C) \delta_i \mu \right] = 0 \tag{0.3}$$

on ω_h (the summation from 1 to 3 is assumed over the repeated indices i), with

$$J_{lS} = -(s_l(\tau_0 \rho)) \left[\delta_l(G_h - \Phi) - (s_l \mu) \delta_l C \right], \ l = 1, 2, 3, \tag{0.4}$$

Copyright © 2021 The Author(s). Published by Vilnius Gediminas Technical University This is an Open Access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided the original author and source are credited.

Recall that $\tau_0 = \tau(\rho, \mathbf{u}, C)|_{\mathbf{u}=\mathbf{0}} > 0, \ \Phi \in H_X(\omega_h)$ is a given function and

$$G_{h} = \Psi_{1\rho}'(\rho, C) + \frac{1}{2}\lambda_{1}s_{i}^{*}\left[(\delta_{i}C)^{2}\right], \quad \mu = \frac{1}{\rho}\left[\Psi_{1C}'(\rho, C) - \delta_{i}^{*}\left(\lambda_{1}(s_{i}\rho)\delta_{i}C\right)\right] \quad (0.5)$$

with the partial derivatives of a given function $\Psi_1(\rho, C)$.

Theorem 2. The equilibrium solutions satisfy the following equations

$$\Psi_{1\rho}'(\rho,C) + \frac{1}{2}\lambda_1 s_i^* \left[(\delta_i C)^2 \right] - \mu_S C - \Phi \equiv G_h - \mu_S C - \Phi \equiv const, \qquad (0.6)$$

$$\Psi_{1C}'(\rho, C) - \delta_i^* \left(\lambda_1(s_i \rho) \delta_i C \right) = \mu_S \rho, \quad \mu_S \equiv const \tag{0.7}$$

on ω_h , with the same functions Ψ_1 and Φ as in the differential case, see (2.12)–(2.13) in [1].

Proof. We first apply the known formula (for example, see formula (14) in [3])

$$\delta_l^*(J_{lS}s_lC) = (\delta_l^*J_{lS})C + s_l^*(J_{lS}\delta_lC)$$

and equation (0.1) and rewrite equation (0.3) as

$$s_i^*(J_{iS}\delta_i C) - \delta_i^*[M(s_i C)\delta_i \mu] = 0 \quad \text{on} \ \omega_h.$$

$$(0.8)$$

We take the inner product in $H_X(\omega_h)$ of equations (0.1) and (0.8) respectively by $G_h - \Phi$ and μ and add the results. We apply both formulas (3.1) in [1] and get

$$-(J_{iS},\delta_i(G_h-\Phi))_{i*} + (J_{iS}\delta_iC,s_i\mu)_{i*} + (M(s_iC),(\delta_i\mu)^2)_{i*} = 0.$$
(0.9)

The substitution of expression (0.4) into (0.9) leads to the equality

$$(s_i(\tau_0\rho), [\delta_i(G_h - \Phi) - (s_i\mu)\delta_iC]^2)_{i*} + (M(s_iC), (\delta_i\mu)^2)_{i*} = 0.$$

Since $\tau_0 \rho > 0$ and $M(s_i C) > 0$, it leads to the equalities

$$\delta_i(G_h - \Phi) - (s_i\mu)\delta_i C = 0, \ \delta_i\mu = 0 \ \text{on} \ \omega_{i*,h}, \ i = 1, 2, 3.$$

This first implies that $\mu \equiv \mu_S = \text{const}$ and then $G_h - \Phi - \mu_S C \equiv \text{const}$ on ω_h .

Consequently $J_{lS} = 0$, l = 1, 2, 3, and all the equations (0.1)–(0.4) are reduced to the two found constancy properties, i.e., according to definitions (0.5), to system (0.6)–(0.7) for ρ_S and C_S on ω_h .

Notice that above we have not required for the regularization parameter τ_0 to be constant as in [1].

References

- V. Balashov and A. Zlotnik. An energy dissipative spatial discretization for the regularized compressible Navier-Stokes-Cahn-Hilliard system of equations. *Math. Model. Anal.*, 25(1):110–129, 2020. https://doi.org/10.3846/mma.2020.10577.
- [2] V. Balashov and A. Zlotnik. On a new spatial discretization for a regularized 3d compressible isothermal Navier-Stokes-Cahn-Hilliard system of equations with boundary conditions. J. Sci. Comput., 86, 2021. https://doi.org/10.1007/s10915-020-01388-6.
- [3] A.A. Zlotnik. Spatial discretization of the one-dimensional quasi-gasdynamic system of equations and the entropy balance equation. *Comput. Math. Math. Phys.*, 52(7):1060–1071, 2012. https://doi.org/10.1134/S0965542512070111.